# A Recursive Inequality of Empirical Measures Associated with EDM 

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#### Abstract

Summary In this paper we consider a random model related to stochastic interacting systems, named an environment-dependent spatial model (EDM). As a matter of fact, this stochastic model is deeply connected with some Markov processes investigated by Liggett [18]. We shall show that rescaled processes of the empirical measures derived from EDMs satisfy some applicationally important recursive inequality.


Key Words: environment-dependent model, random model, stochastic interacting systems, empirical measures, rescaled process, recursive inequality.

## 1. Introduction

In this section we shall introduce an environment-dependent random model (EDM)[12]. Let $\mathbb{Z}^{d}$ be a $d$-dimensional lattice space, and we suppose that each site on $\mathbb{Z}^{d}$ is occupied by all means by either one of the two species. Just after a random time period, a particle dies out and is replaced by a new one, but the random time and the type chosen of the species are assumed to be determined by the environment conditions around the particle. The random function $\xi_{t} \equiv \xi_{t}(x): \mathbb{Z}^{d} \rightarrow$ $\{0,1\}$ denotes the state at time $t$, and each number of $\{0,1\}$ denotes the label of the type chosen of the two species. When we set $\|y\|_{\infty}:=\max _{i} y_{i}$ for $y=\left(y_{1}, \ldots, y_{\mathrm{d}}\right)$, then the $R$-neighborhood of $x$ is defined by

$$
\begin{equation*}
\mathcal{N}_{x}:=x+\left\{y: 0<\|y\|_{\infty} \leqslant R\right\}, \tag{1}
\end{equation*}
$$

where $R$ is a positive constant given. For $i=0,1$, let $f_{i}(x, \xi)$ be a frequency of appearance of type $i$ in $\mathcal{N}_{x}$ for $\xi$. More precisely, it can be expressed as

$$
\begin{equation*}
f_{i}(x) \equiv f_{i}(x, \xi):=\frac{\#\left\{y: \xi_{t}(y)=i ; y \in \mathcal{N}_{x}\right\}}{\# \mathcal{N}_{x}} . \tag{2}
\end{equation*}
$$

For non-negative parameters $\alpha_{i j} \geq 0$, the dynamics of $\xi_{t}$ is defined as follows. The state $\xi$ makes transition $0 \rightarrow 1$ at rate

$$
\begin{equation*}
\frac{\lambda f_{1}\left(f_{0}+\alpha_{01} f_{1}\right)}{\lambda f_{1}+f_{0}} \tag{3}
\end{equation*}
$$

and it makes transition $1 \rightarrow 0$ at rate

$$
\begin{equation*}
\frac{f_{0}\left(f_{1}+\alpha_{10} f_{0}\right)}{\lambda f_{1}+f_{0}} . \tag{4}
\end{equation*}
$$

The interpretation of the above rate is as follows. The particle of type $i$ dies at rate $f_{i}+\alpha_{i j} f_{j}$, and is replaced instantaneously by either one of the two species chosen at random, according to the proliferation rate of type 0 and the interaction (= the competitive result) with the particle of type 1. We assume that competitive two species possess the same intensity of intraspecific interaction. The exchange of particles after death is described in the form being proportional to the weighted density between the two species, expressed by a parameter $\lambda$. Assume usually that $\lambda \geq 1$. The case of $\lambda$ $=1$ means that the contribution to a local appearance rate between the two competitive species is equivalent.

## 2. Scaling, rescaled process and empirical measure

For simplicity we shall treat a simple case $\lambda=1$ only in what follows. For $N=1,2, \ldots$, let $M_{N} \in \mathbb{N}$, and we put $\ell_{N}:=M_{N} \sqrt{N}$, and $\mathbb{S}_{N}:=\mathbb{Z}^{d} / \ell_{N}$. And also $W_{N}=\left(W_{N}^{1}, \ldots, W_{N}^{d}\right) \in\left(\mathbb{Z}^{d} / M_{N}\right) \backslash$ $\{0\}$ is defined as a random vector satisfying
(i) $\mathcal{L}\left(W_{N}\right)=\mathcal{L}\left(-W_{N}\right)$;
(ii) $E\left(W_{N}^{i} W_{N}^{j}\right) \rightarrow \delta_{i j} \sigma^{2}(\geq 0) \quad($ as $N \rightarrow \infty)$;
(iii) $\left\{\left|W_{N}\right|^{2}\right\}(N \in \mathbb{N})$ is uniformly integrable.

Here $\mathcal{L}(Y)$ indicates the law of a random variable $Y$. For the kernel $p_{N}(x):=P\left(W_{N} / \sqrt{N}=x\right), x \in$ $\mathbb{S}_{N}$ and $\xi \in\{0,1\}^{\mathbb{S}_{N}}$, we define the scaled frequency $f_{i}^{N}$ as

$$
\begin{equation*}
f_{i}^{N}(x, \xi)=\sum_{y \in \mathbb{S}_{N}} p_{N}(y-x) 1_{\{\xi(y)=i\}}, \quad(i=0,1) . \tag{5}
\end{equation*}
$$

Actually, $\xi_{t}^{N}$ is given by $\xi_{t}^{N}=\xi_{N t}(\mathrm{x} \sqrt{N})$. As a matter of fact, the rescaled process $\xi_{t}^{N}: \mathbb{S}_{N} \ni x \mapsto \xi_{t}^{N}$ $(x) \in\{0,1\}$ is determined by the following state transition law, nemaly, it makes transition $0 \rightarrow 1$ at rate $N f_{1}^{N}\left(\mathrm{f}_{0}^{N}+\alpha_{0}^{N} f_{1}^{N}\right)$, or else it makes transition $1 \rightarrow 0$ at rate $N f_{0}^{N}\left(f_{1}^{N}+\alpha_{1}^{N} f_{0}^{N}\right)$. We also denote the rescaled process $\xi_{t}^{N}$ by the symbol $\operatorname{Res}\left(p_{N}, \alpha_{i}^{N}\right)$. On this account, we may define the associated measure-valued process (or its corresponding empirical measure) as

$$
\begin{equation*}
X_{t}^{N}:=\frac{1}{N} \sum_{x \in \mathbb{S}_{N}} \xi_{t}^{N}(x) \delta_{x} . \tag{6}
\end{equation*}
$$

For the initial value $X_{0}^{N}$, we assume that

$$
\begin{equation*}
\sup _{N}\left\langle X_{0}^{N}, 1\right\rangle<\infty, \quad X_{0}^{N} \rightarrow X_{0} \quad \text { in } \quad M_{F}\left(\mathbb{R}^{d}\right) \quad(N \rightarrow \infty), \tag{7}
\end{equation*}
$$

where $M_{F}\left(\mathbb{R}^{d}\right)$ is the totality of all the finite measures on $\mathbb{R}^{d}$, equipped with the topology of weak convergence. For a finite measure $\mu \in M_{F}(E)$ with a topological space $E$, we use the notation $\langle\mu$, $\varphi\rangle=\int_{E} \varphi(x) \mu(d x)$ for integral of a measurable function $\varphi$ over $E$ with respect to a measure $\mu$ on $E$. Note that the convergence in (7) is that in the sense of weak convergence for measures [17].

## 3. Main theorem : recursive inequality

In this section we shall introduce the principal result on an estimate of the maximum of the moment of total mass process for the empirical measure. To prove it we need some precise esti-
mate of the quantity in question, and in fact, that can be realized by a certain recursive type inequality for the empirical measures.

Theorem 1. (Main Result) Let $F(N)$ be a function of $N$ that satisfies $1 \leqslant F(N) \leqslant N$ and $\lim _{N \rightarrow \infty} F(N) / N=0$. If the condition $N^{5 / 7} / F(N) \rightarrow 0$ holds as $N \rightarrow \infty$, then for any $p>1$ and $T$ $>0$, there exists a finite constant $c(p, T)>0$ such that

$$
\begin{equation*}
E\left[\sup _{t \leqslant T}\left\langle X_{t}^{N}, 1\right\rangle^{p}\right] \leqslant c(p, T)\left(\frac{N}{F(N)}\right)^{p-1 / 2}\left(\left\langle X_{0}^{N}, 1\right\rangle^{p}+1\right) . \tag{8}
\end{equation*}
$$

## 4. Sketch of proof of main result

Step 1. In this section we shall introduce a sketch of proof of our main result Theorem 1. First of all, we begin with showing a useful equality.

Lemma 2. The following equality holds for every $t>0$ :

$$
\begin{equation*}
E\left[\left\langle X_{t}^{N}, 1\right\rangle\right]=\left\langle X_{0}^{N}, 1\right\rangle . \tag{9}
\end{equation*}
$$

Proof. First we consider a bounded function $\psi: \mathbb{S}_{N} \rightarrow \mathbb{R}$. For a continuous time random walk $B_{t}^{x_{N} N}$ with rate $N$ and step distribution $p_{N}$ starting at $x$,

$$
\begin{equation*}
\phi_{s}(x) \equiv \phi(s, x)=P_{t-s}^{N} \psi(x):=E\left[\psi\left(B_{t-s}^{x, N}\right)\right] \tag{10}
\end{equation*}
$$

defines a semigroup. Indeed, this newly defined function $\phi$ satisfies a differential equation $\partial_{s} \phi(s)$ $+\mathcal{A}_{N} \phi(s)=0$ by virtue of the backward equation argument for continuous time Markov chains. Recall that $\mathcal{A}_{N}$ is its generator, and is given by

$$
\begin{equation*}
\mathcal{A}_{N} \phi(x):=N \sum_{y} p_{N}(y-x)(\phi(y)-\phi(x)) . \tag{11}
\end{equation*}
$$

According to the theory of semimartingales [21], we may apply Itô's formula in stochastic calculus [14] to a relationship of rescaled EDMs to obtain

$$
\begin{equation*}
\left\langle X_{t}^{N}, \phi_{t}\right\rangle=\left\langle X_{0}^{N}, \phi_{0}\right\rangle+\int_{0}^{t} X_{s}^{N}\left(\partial_{s} \phi(s)+\mathcal{A}_{N} \phi(s)\right) d s+M_{t}^{N}(\phi), \tag{12}
\end{equation*}
$$

for $0 \leqslant t \leqslant T$, where $M_{t}^{N}(\phi)$ is a martingale term and $X_{t}(\phi)$ denotes an integral of the test function $\phi$ relative to the measure-valued process $d X_{t}$. Taking the expectation operation $E[\cdot]$ at the both sides of (12), we can get the equality $E\left[X_{t}^{N}\left(\phi_{t}\right)\right]=E\left[X_{0}^{N}\left(\phi_{0}\right)\right]=E\left[X_{0}^{N}\left(P_{t}^{N} \psi\right)\right]$, and besides we readily obtain the desired expression (9) with $E\left[X_{t}^{N}(\psi)\right]=E\left[X_{0}^{N}\left(P_{t}^{N} \psi\right)\right]$ by changing the function $\phi(t)$ to a general one $\psi$, where we have substituted $\phi \cdot(x) \equiv 1$ instead of $\psi$ and we also have made use of

$$
\begin{equation*}
\phi(s, x)=\phi_{s}(x)=P_{t-s}^{N} 1(x)=E\left[1\left(B_{t-s}^{x, N}\right)\right]=E[1(x)]=1 \tag{13}
\end{equation*}
$$

This finishes the proof of lemma.

Step 2. Recall standard results for stochastic integrals with respect to Poisson processes $N_{s}$ with the
intensity $E\left[N_{s}\right]=\lambda_{s}$. Since $\hat{N}_{s}=N_{s}-\lambda_{s}$ is a martingale, the stochastic integral $M_{s}=\int_{0}^{t} \Psi(s, \omega)$ $d \hat{N}_{s}$ becomes a martingale. Furthermore, it follows that

$$
\begin{equation*}
E\left|M_{t}\right|^{2}=E\left|\int_{0}^{t} \Psi(s, \omega) d \hat{N}_{s}\right|^{2}=E \int_{0}^{t} \Psi^{2}(s, \omega) d \lambda_{s} \tag{14}
\end{equation*}
$$

Let $\left\{\Lambda_{t}^{N}(x, y): x, y \in \mathbb{S}_{N}\right\}$ be a family of independent Poisson processes with rate $N \cdot p_{N}(y-x)$ defined on a complete probability space. Note that its compensated process

$$
\begin{equation*}
\hat{\Lambda}_{t}^{N}(x, y)=\Lambda_{t}^{N}(x, y)-N \cdot p_{N}(y-x) t \tag{15}
\end{equation*}
$$

are $\left(\mathcal{F}_{t}\right)$-martingale. On this account, for every test function $\psi \equiv \psi(s, x) \in M_{b}\left([0, T] \times \mathbb{S}_{N}\right)$ with $T$ $<\infty$,

$$
\begin{equation*}
M_{t}^{N}(\psi):=\frac{1}{F(N)} \sum_{s} \sum_{y} \int_{0}^{t} \psi_{s}(x)\left(\xi_{s-}(y)-\xi_{s-}(x)\right) d \hat{\Lambda}_{s}(x, y) \tag{16}
\end{equation*}
$$

is a cadlag $L^{2}\left(\mathcal{F}_{t}\right)$-martingale, and its predictable square function is given by

$$
\begin{equation*}
\left\langle M^{N}(\psi)\right\rangle_{t}=\frac{N}{F(N)^{2}} \int_{0}^{t} \sum_{x} \sum_{y} \psi_{s}(x)^{2}\left(\xi_{s}(y)-\xi_{s}(x)\right)^{2} p_{N}(y-x) d s, \quad t \in[0, T] \tag{17}
\end{equation*}
$$

where $M_{b}(D)$ is the totality of all bounded measurable functions defined on a proper space $D$, and the summation $\sum_{x}$ is taken over the whole space $\mathbb{S}_{N}$. In particular, the equality

$$
\begin{equation*}
\left\langle M^{N}(1)\right\rangle_{t}=2 \int_{0}^{t}\left\langle X_{s}^{N}, \frac{N}{F(N)} V_{N}(s, x)\right\rangle d s \tag{18}
\end{equation*}
$$

holds, where $V_{N}(t, x)=\sum_{y} p_{N}(y-x) 1\left\{\xi_{t}(y)=0\right\}$. An application of the results (16), (17) and (18) with the expression (12) leads to $X_{t}^{N}(1)=X_{0}^{N}(1)+\mathrm{M}_{t}^{N}(1)$.

Lemma 3. The random quantity $\left\langle X_{t}^{N}, 1\right\rangle$ is an $L^{2}$-martingale such that

$$
\begin{equation*}
\left\langle X^{N}(1)\right\rangle_{t}=\frac{2 N}{F(N)} \int_{0}^{t} \frac{1}{F(N)} \sum_{x} \xi_{s}(x) V_{N}(s, x) d s \leqslant \frac{2 N}{F(N)} \int_{0}^{t}\left\langle X_{s}^{N}, 1\right\rangle d s \tag{19}
\end{equation*}
$$

holds.
Proof. The expression (16) implies that $\mathrm{M}_{t}^{N}(1)$ is also a martingale as a special case. While, since $\left\langle X_{0}^{N}, 1\right\rangle=$ const., it is easy to see that for $0<\forall s<t$

$$
\begin{align*}
E\left[X_{t}^{N}(1) \mid \mathcal{F}_{s}\right] & =E\left[X_{0}^{N}(1) \mid \mathcal{F}_{s}\right]+E\left[M_{t}^{N}(1) \mid \mathcal{F}_{s}\right]  \tag{20}\\
& =X_{0}^{N}(1)+M_{s}^{N}(1)=X_{s}^{N}(1) .
\end{align*}
$$

Moreover, we can deduce that $\left\langle X_{t}^{N}, 1\right\rangle$ is an $L^{2}$ martingale because $M_{t}^{N}(\psi)$ is an $L^{2}$ martingale. For a proper sequence of stopping times $\left(T_{n}\right)$, the quadratic variation satisfies

$$
\begin{equation*}
\operatorname{Var}\left(\left\langle X_{0}^{N}, 1\right\rangle\right):=\sum_{k}\left|\left\langle X^{N}\left(0 \wedge T_{k+1}\right), 1\right\rangle-\left\langle X^{N}\left(0 \wedge T_{k}\right), 1\right\rangle\right|^{2}=0, \tag{21}
\end{equation*}
$$

hence it follows immediately that $\left\langle X^{N}(1)\right\rangle_{t}=\left\langle M_{t}^{N}(1)\right\rangle_{t}$ holds for every $t>0$. Elementary results for Poisson process and stochastic integral with respect to Poisson process reads

$$
\begin{align*}
E\left|M_{t}^{N}(1)\right|^{2} & =E\left|\frac{1}{F(N)} \sum_{x} \sum_{y} \int_{0}^{t}\left(\xi_{s-}(y)-\xi_{s-}(x)\right) d \hat{\Lambda}_{s}(x, y)\right|^{2}  \tag{22}\\
& =E\left[\frac{1}{F(N)^{2}} \sum_{x} \sum_{y} \int_{0}^{t}\left(\xi_{s}(y)-\xi_{s}(x)\right)^{2} N \cdot p_{N}(y-x) d s\right]
\end{align*}
$$

From this fact, we can get easily that

$$
\begin{align*}
\left\langle M_{t}^{N}(1)\right\rangle_{t}= & \frac{N}{F(N)^{2}} \sum_{x} \sum_{y} \int_{0}^{t}\left(\xi_{s}(y)-\xi_{s}(x)\right)^{2} p_{N}(y-x) d s \\
= & \frac{N}{F(N)^{2}} \sum_{x} \int_{0}^{t} \xi_{s}(x)\left(\sum_{y} p_{N}(y-x) \cdot 1\left\{\xi_{s}(x)=0\right\}\right) d s \\
& \quad+\frac{N}{F(N)^{2}} \sum_{y} \int_{0}^{t} \xi_{s}(y)\left(\sum_{x} p_{N}(x-y) \cdot 1\left\{\xi_{s}(x)=0\right\}\right) d s  \tag{23}\\
= & \frac{2 N}{F(N)} \int_{0}^{t}\left(\frac{1}{F(N)} \sum_{x} \xi_{s}(x) V_{N}(s, x)\right) d s \\
\leqslant & \frac{2 N}{F(N)} \int_{0}^{t}\left(\frac{1}{F(N)} \sum_{x} \xi_{s}(x) 1(x)\right) d s=\frac{2 N}{F(N)} \int_{0}^{t}\left\langle X_{s}^{N}, 1\right\rangle d s
\end{align*}
$$

where we made use of a simple inequality $V_{N}(s, x) \leqslant 1$.
Step 3. To complete the proof of Theorem 1, we need the following lemmas.
Lemma 4. There exists a positive constant $K(p)>0$ depending on $\forall p>1$ such that

$$
\begin{equation*}
E\left[\sup _{t \leqslant T}\left\langle X_{t}^{N}, 1\right\rangle^{p}\right] \leqslant K(p)\left\{X_{0}^{N}(1)^{p}+E\left[\left\langle X^{N}(1)\right\rangle_{t}^{p / 2}+E\left[\sup _{t \leqslant T}\left|\Delta X_{t}^{N}(1)\right|^{p}\right]\right\}\right. \tag{24}
\end{equation*}
$$

holds.
Lemma 5. There exists a positive constant $K_{1}(p)>0$ such that

$$
\begin{equation*}
E\left[\sup _{t \leqslant T}\left\langle X_{t}^{N}, 1\right\rangle^{p}\right] \leqslant K_{1}(p)\left\{X_{0}^{N}(1)^{p}+\left(\frac{2 N T}{F(N)}\right)^{p / 2} E\left[\sup _{t \leqslant T} X_{t}^{N}(1)^{p / 2}\right]+1\right\} \tag{25}
\end{equation*}
$$

holds.
Lemma 6. There exists a positive constant $K_{2}>0$ such that

$$
\begin{equation*}
E\left[\sup _{t \leqslant T} X_{t}^{N}(1)^{2}\right] \leqslant K_{2}\left(\frac{N T}{F(N)}+1\right)\left(X_{0}^{N}(1)^{2}+1\right) \tag{26}
\end{equation*}
$$

holds.
Lemme 7. There exist certain positive constants $K_{3}(p)>0$ and $c(p, T)>0$ such that

$$
\begin{equation*}
E\left[\sup _{t \leqslant T}\left\langle X_{t}^{N}, 1\right\rangle^{p}\right] \leqslant K_{3}(p)\left[X_{0}^{N}(1)^{p}+\left(\frac{2 N t}{F(N)}\right)^{p / 2}\left\{c(p, T)\left(\frac{N}{F(N)}\right)^{(p-1) / 2}\left(X_{0}^{N}(1)+1\right)\right\}+1\right] \tag{27}
\end{equation*}
$$

holds.
The above-mentioned Lemma 4 is a direct result of Burkholder-Davis-Gundy inequality [21] together with a trivial inequality $(a+b)^{p} \leqslant c_{p}\left(a^{p}+b^{p}\right)$ and $\langle X\rangle_{t}=\langle X\rangle_{\mathrm{t}}^{c}+\sum_{0 \leqslant s \leqslant \mathrm{t}}\left(\Delta X_{s}\right)^{2}$ and $\Delta X_{s}=$ $\Delta M_{s}$ for $\forall_{s}$. Lemma 5 yields from a standard technique for estimation, the inequality (19) in Lemma 3, and an estimate for jump term:

$$
\begin{equation*}
E\left[\sup _{s \leqslant T}\left|\Delta X_{s}^{N}(1)\right|^{p}\right] \leqslant\left(\frac{1}{F(N)}\right)^{p} \leqslant 1 . \tag{28}
\end{equation*}
$$

Another trivial inequality $a b \leqslant \frac{1}{2}\left(a^{2}+b^{2}\right)$ yields to Lemma 6 . Moreover, Lemma 7 can be verified by induction hypothesis. Combining those results, we can deduce that (8) is valid for $\forall p>1$ by taking advantage of mathematical induction method. This completes the proof of Theorem 1.

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## References

1. Cox, J. T., Durrett, R. and Perkins, E. A. : Rescaled voter models converge to super- Brownian motion. Ann. Probab. 28 (2000), 185-234.
2. Dôku, I. : Exponential moments of solutions for nonlinear equations with catalytic noise and large deviation. Acta Appl. Math. 63 (2000), 101-117.
3. Dôku, I. : Weighted additive functionals and a class of measure-valued Markov processes with singular branching rate. Far East J. Theo. Stat. 9 (2003), 1-80.
4. Dôku, I. : A certain class of immigration superprocesses and its limit theorem. Adv. Appl. Stat. 6 (2006), 145-205.
5. Dôku, I. : A limit theorem of superprocesses with non-vanishing deterministic immigration. Sci. Math. Japn. 64 (2006), 563-579.
6. Dôku, I. : Limit theorems for rescaled immigration superprocesses. RIMS Kôkyûroku Bessatsu, B6 (2008), 56-69.
7. Dôku, I. : A limit theorem of homogeneous superprocesses with spatially dependent parameters. Far East J. Math. Sci. 38 (2010), 1-38.
8. Dôku, I. : On mathematical modeling for immune response to the cancer cells. J. SUFE Math. Nat. Sci. 60 (2011), no.1, 137-148.
9. Dôku, I. : Limit Theorems for Superprocesses: Rescaled Processes, Immigration Superprocesses and Homogeneous Superprocesses. Lap Schalt. Lange, Berlin, 2014.
10. Dôku, I. : Star-product functional and unbiased estimator of solutions to nonlinear integral equations. Far East J. Math. Sci. 89 (2014), 69-128.
11. Dôku, I. : Some integral equation related to a branching model. RIMS Kôkyûroku (Kyoto Univ.) 1937 (2015), 25-31.
12. Dôku, I. : A example for convergence of environment-dependent spatial models. J. SUFE Math. Nat. Sci. 65 (2016), no.1, 179-186.
13. Dôku, I. : Tumour immunoreaction and environment-dependent models. To appear in Trans. Japn. Soc. Indu. Appl. Math. (2016), pp.1-40.
14. Durrett, R. : Stochastic Calculus: A Practical Introduction. CRC Press, Boca Raton, 1996.
15. Ikeda, N. and Watanabe, S. : Stochastic Differential Equations and Diffusion Processes. Second Edition, North-Holland, Amsterdam, 1989.
16. Jacob, J. and Shiryaev, A. N. : Limit Theorems for Stochastic Processes. A series of Comprehensive Studies in Mathematics 288, Springer, Berlin, 1987.
17. Kallenberg, O. : Foundations of Modern Probability. Second Edition, Springer, New York, 2002.
18. Liggett, T. M. : Interacting Particle Systems. Springer, New York, 1985.
19. Liggett, T. M. : Stochastic Interacting Systems: Contact, Voter and Exclusion Processes. Springer, New York, 1999.
20. Liggett, T. M. : Continuous Time Markov Processes: An Introduction. Graduate Studies in Math. vol. 113, Amer. Math. Soc. Providence, 2010.
21. Protter, P. E. : Stochastic Integration and Differential Equations. Second Edition, A Series in Stochastic Modelling and Applied Probability 21, Springer, Heidelberg, 2004.

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